

25 XI 2013

Procedura grupy normalizacji i jej znaczenie w fizyce

(1)

$$W = 2^{N_A}$$



$$E = \sum_{i=1}^N S_i S_{i+1} \cdot E_0$$

$$Z = \sum_{\{S_i\}_{i=1}^N} e^{-E/kT}$$

$$e^{-iHt/\hbar}$$

↑
E

$$H = M_{2 \times 2} = \begin{pmatrix} a & b \\ b & c \end{pmatrix}$$

$$M_{2 \times 2} \left(x \begin{bmatrix} 1 \\ 0 \end{bmatrix} + y \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right) \rightarrow \begin{bmatrix} ax + by \\ bx + cy \end{bmatrix}$$

$$S_{i+1} = M_{2 \times 2} S_i$$

Fizyka statystyczna (klasyka) i Mechanika Kwantowa są powiązane ze sobą

$$MK \quad \frac{d+l}{v \quad t} = D \quad \leftrightarrow \quad FS \quad \frac{D}{v}$$

Wzrosty jest dwa

$$|1\rangle \quad |2\rangle$$

Stany

$$H|1\rangle = a|1\rangle + b|2\rangle$$

$$H|2\rangle = c|2\rangle + b|1\rangle$$

$$\langle 1|H|1\rangle = a \langle 1|1\rangle = a$$

$$\langle 2|H|2\rangle = c$$

$$\langle 2|H|1\rangle = b = \langle 1|H|2\rangle$$

$$e^{-iHt/\hbar}$$

$$e^{-iH\frac{t}{\hbar}} e^{-iH\frac{t}{\hbar}} \dots$$

$$e^{-E/kT} = e^{-\sum E_i/kT}$$

$$= e^{-E_1/kT} e^{-E_2/kT} \dots$$

$$= e^{-E_1/kT} e^{-E_2/kT} \dots$$

(2.)

$$M_{2 \times 2} = H_2 = \begin{bmatrix} a & b \\ b & c \end{bmatrix}$$

$$2 \rightarrow 3$$

$$M_{3 \times 3} = H_3 = \begin{bmatrix} d & e & f \\ e & g & h \\ f & h & i \end{bmatrix}$$

$$L_p \sim 10^{80}$$

$$3 \rightarrow 2^{10^{80}} ?$$

~ linie protand
w naszym
wszedliście

$$b > 1, N \times N, H_{mn} = \langle m | H | n \rangle = b^n \delta_{mn} - g b^{\frac{m+n}{2}} \left\{ \begin{array}{l} \text{kandydat} \\ \text{na} \\ \text{prawo} \\ \text{pryrody} \end{array} \right. \quad (3)$$

$$E_1 = b, E_2 = b^2, E_3 = b^3, E_{-16} = b^{-16}$$

$$E \sim 10^{10} - E \sim 10^{-10}$$

zakres
21 x 21

$$H_{mn} = E_m \delta_{mn} - g \sqrt{E_m E_n}$$

Jakże stany własne w spójnej spójności (praktycznie stabilny),

$$e^{-iHt/\hbar} \psi, H\psi_\lambda = \lambda \psi_\lambda, e^{-iHt/\hbar} \psi_\lambda = e^{-i(E_\lambda = \lambda)t/\hbar}$$

$$H_{mn} = b^m \delta_{mn} - g b^{\frac{m+n}{2}}$$

Newton \rightarrow George E. Smith

Stanford U., Lectures

"The Missing Link" ?

(4)

$$\sum_{n=1}^{M+1} H_{mn} \psi_n = \lambda \psi_m \quad m \in [M, N]$$

$M \geq 0 \quad N > 0$

$$m \quad b^m \psi_m - g \sum_{n=M}^N b^{\frac{m+n}{2}} \psi_n = \lambda \psi_m, \quad N-M+1 \text{ remain, all unspecified } m.$$

$$b^m \psi_m - g b^{\frac{m}{2}} \left(\sum_n b^{\frac{n}{2}} \psi_n \right) = \lambda \psi_m.$$

$$(b^m - \lambda) \psi_m = g b^{\frac{m}{2}} \cdot c, \quad \psi_m = \frac{g c b^{\frac{m}{2}}}{b^m - \lambda}, \quad c = \sum_n b^{\frac{n}{2}} \psi_n = \sum_n \frac{b^{\frac{n}{2}} g b^{\frac{n}{2}} c}{b^n - \lambda}.$$

5.

$$c = \sum_n b^{\frac{n}{2}} \psi_n = \sum_n \frac{b^{\frac{n}{2}} g b^{\frac{n}{2}} c}{b^n - \lambda}$$

$$1 = g \sum_{n=M}^N \frac{b^n}{b^n - \lambda}, \quad b > 1$$

$$1 = g \cos(\lambda) + g(N-k)$$

$N \rightarrow \infty$

$$E_N = b^N \quad \ln E_N = N \ln b = N$$

$b = e$

$$1 \approx g \sum_{n=M}^N \frac{b^n}{- \lambda} + g \sum_{k/b^k \lambda}^N \frac{b^n}{b^n}$$

possibilities
 $b^n < \lambda$

possibilities
 $\lambda < b^n$

$$1 = g \cos(\lambda) + g \left(\ln E_N / E_0 - \ln E_k / E_0 \right)$$

$$1 = g \cos(\lambda) + g \cos, \quad \boxed{g \rightarrow g_N}$$

$$g(\cos \text{ shchitaniya}) + g(N-k) = 1$$

E_λ

$$b^k \approx \lambda = E_\lambda$$

$$I = g \cos(\lambda) + g \ln \frac{E_U}{E_0}$$

$$= g \cos(E_{dosw}) + g \ln \frac{E_U}{E_{dosw}}$$

$\lambda = \frac{E_U}{E_0} = E_{dosw}$

6.

$$I = g \left(\cos + \ln \frac{E_U}{E_{dosw}} \right)$$

$$g_N = \frac{I}{\cos(E_{dosw}) + \ln \frac{E_U}{E_{dosw}}}$$

$$g \rightarrow g(E_U) = g_N$$

Jak to zrozumieć?

Wyjaśnić, gdy powstanie potrzeba.

$$E_U \rightarrow \infty$$

$$g_N \rightarrow 0$$

Asymptotycznie

Stochada.