

Asymptotic freedom of gluons in the Fock space

Stanisław D. Głazek

Faculty of Physics, University of Warsaw

Asymptotic freedom of gluons is described in terms of a family of scale-dependent renormalized Hamiltonian operators acting in the Fock space. The Hamiltonians are obtained by applying the renormalization group procedure for effective particles (RG-PEP) to quantum $SU(3)$ Yang-Mills theory. The RGPEP is a general method for solving quantum field theories in the Minkowski space-time.

S.D.Głazek, *Dynamics of effective gluons*, Phys. Rev. D **63**, 116006 (2001).

S.D.Głazek, *Perturbative formulae for relativistic interactions of effective particles*, Acta Phys. Pol. B43, 1843, 20p (2012).

M.Gómez-Rocha, S.D.Głazek, *Asymptotic freedom in the front-form Hamiltonian for quantum chromodynamics of gluons*, arXiv:1505.06688.

Outline:

explanation of asymptotic freedom of gluons in the Fock space

as an example of application of

the renormalization group procedure for effective particles (**RGPEP**)

as a new computational tool in quantum field theory

In popular approaches to QCD, one way or the other,
one assumes that there exists some H_{QCD} .

$$H_{QCD} = ?$$

simplify

$$H_{YM} = ?$$

$$\begin{aligned}\mathcal{L} &= -\frac{1}{2} \text{tr} F^{\mu\nu} F_{\mu\nu} \\ F^{\mu\nu} &= \partial^\mu A^\nu - \partial^\nu A^\mu + ig[A^\mu, A^\nu] \\ \mathcal{T}^{\mu\nu} &= -F^{a\mu\alpha} \partial^\nu A_\alpha^a + g^{\mu\nu} F^{a\alpha\beta} F_{\alpha\beta}^a / 4 \\ P^\nu &= \int d\sigma_\mu \mathcal{T}^{\mu\nu}\end{aligned}$$

$P^\nu = \int_\Sigma d\sigma_\mu \mathcal{T}^{\mu\nu}$ requires a choice of a hyperplane Σ

$$d\sigma^\mu = (1, 0, 0, 0)^\mu d^3x \quad \mathcal{T}^{00} \quad H = P^0$$

$$d\sigma^\mu = (1, 0, 0, -1)^\mu \frac{1}{2} dx^- d^2x^\perp \quad \mathcal{T}^{+-} \quad H = P^-$$

$$A^\pm = A^0 \pm A^3 \quad A^\perp = (A^1, A^2)$$

instant hyperplane \leftrightarrow 6-dimensional invariance group

front hyperplane \leftrightarrow 7-dimensional invariance group

$$\mathcal{H}_{YM} = \mathcal{H}_{A^2} + \mathcal{H}_{A^3} + \mathcal{H}_{A^4}$$

$$\mathcal{H}_{A^3} = g i \partial_\alpha A_\beta^a [A^\alpha, A^\beta]^a$$

quantization $\hat{A}^\mu = \sum_{\sigma c} \int_k \left[t^c \varepsilon_{k\sigma}^\mu a_{k\sigma c} e^{-ikx} + t^c \varepsilon_{k\sigma}^{\mu*} a_{k\sigma c}^\dagger e^{ikx} \right]_{on \Sigma}$

$$\left[a_{k\sigma c}, a_{k'\sigma'c'}^\dagger \right] = k^+ \tilde{\delta}(k - k') \delta^{\sigma\sigma'} \delta^{cc'}$$

$$\hat{H}_{YM} = \frac{1}{2} \int dx^- d^2 x^\perp : \mathcal{H}_{YM}(\hat{A}) :$$

Is it it?

No, it is not.

Key example

$$\hat{H}_{A^3} = \int_{\Sigma} g : i\partial_{\alpha}\hat{A}_{\beta}^a[\hat{A}^{\alpha}, \hat{A}^{\beta}]^a :$$

$$\hat{A} \sim a + a^{\dagger} \quad : \hat{A}^3 : \sim a^{\dagger 3} + a^{\dagger 2}a + a^{\dagger}a^2 + a^3$$

$$a_k|0\rangle = 0 \quad a_{k_1}^{\dagger} a_{k_2}^{\dagger} a_{k_3}^{\dagger} |0\rangle = ? \quad e^{-i\hat{H}t/\hbar}|0\rangle = ?$$

$$\Delta E_{|0\rangle}^{(2)} = \sum_{|3g\rangle} \frac{|\langle 3g|\hat{H}|0\rangle|^2}{-E_{3g}} = -\infty$$

$$\Delta E_{|g\rangle}^{(2)} = \sum_{|2g\rangle} \frac{|\langle 2g|\hat{H}|g\rangle|^2}{E_g - E_{2g}} = -\infty \quad \text{etc.}$$

$$\hat{H}|G P\rangle = \frac{M_G^2 + P^{\perp 2}}{P^+} |G P\rangle$$

$$P^2 = P^+ P^- - P^{\perp 2} = M_G^2$$

$$|G P\rangle = |gg P\rangle + |ggg P\rangle + |gggg P\rangle + \dots$$

$$(P^+ \hat{H} - P^{\perp 2})|G P\rangle = M_G^2 |G P\rangle$$

infinitely large range of momenta (locality)

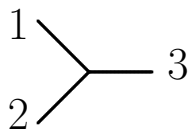
infinitely many components (products of more than 2 fields)

Regularization

example $\hat{H}_{A^3} \rightarrow \hat{H}_{A^3 R}$

$$\hat{H}_{A^3} = \sum_{123} \int [123] \delta(1 + 2 - 3) \left[g Y_{123} a_1^\dagger a_2^\dagger a_3 + g Y_{123}^* a_3^\dagger a_2 a_1 \right]$$

$$\rightarrow \hat{H}_{A^3 R} = \sum_{123} \int [123] \delta(1 + 2 - 3) R \left[g Y_{123} a_1^\dagger a_2^\dagger a_3 + g Y_{123}^* a_3^\dagger a_2 a_1 \right]$$



$$x_1 = p_1^+ / p_3^+ = x$$

$$k_1 = p_1 - x_1 p_3 = \kappa$$

$$x_2 = p_2^+ / p_3^+ = 1 - x$$

$$k_2 = p_2 - x_2 p_3 = -\kappa$$

$$r(x_i, k_i^\perp) = x_i^\delta e^{-k_i^{\perp 2} / \Delta^2} \quad i = 1, 2, 3 \quad x_3 = 1 \quad k_3^\perp = 0^\perp$$

$$R = r(x_1, k_1^\perp) r(x_2, k_2^\perp) r(1, 0^\perp) = x^\delta (1 - x)^\delta e^{-2\kappa^2 / \Delta^2}$$

$$Y_{123} = i f^{c_1 c_2 c_3} \left[\varepsilon_1^* \varepsilon_2^* \cdot \varepsilon_3 \kappa - \varepsilon_1^* \varepsilon_3 \cdot \varepsilon_2^* \kappa \frac{1}{x_2} - \varepsilon_2^* \varepsilon_3 \cdot \varepsilon_1^* \kappa \frac{1}{x_1} \right]$$

$$\hat{H}_R |G P\rangle = \frac{M_G^2 + P^\perp{}^2}{P^+} |G P\rangle$$

$$|G P\rangle = |gg P\rangle + |ggg P\rangle + |gggg P\rangle + \dots$$

$$(P^+ \hat{H}_R - P^\perp{}^2) |G P\rangle = M_G^2 |G P\rangle$$

momenta limited by the regularization

need to remove effects of regularization \rightarrow

renormalization of Hamiltonians \rightarrow

scale-dependent effective theory, including bound states

$$\begin{array}{c} \text{bare gluons of size 0} \\ \left[\begin{array}{c} \dots \\ |gggggg\rangle \\ |ggggg\rangle \\ |gggg\rangle \\ |ggg\rangle \\ |gg\rangle \end{array} \right] \end{array} = \begin{array}{c} \text{gluons of size } s \\ \left[\begin{array}{c} \dots \\ |ggggg\rangle \\ |gggg\rangle \\ |ggg\rangle \\ |gg\rangle \\ |g\rangle \end{array} \right] \otimes \left[\begin{array}{c} \dots \\ |ggggg\rangle \\ |gggg\rangle \\ |ggg\rangle \\ |gg\rangle \\ |g\rangle \end{array} \right] + \dots
 \end{array}$$

$$|gg\rangle + |ggg\rangle + \dots = |g_s g_s\rangle + |g_s g_s g_s\rangle + \dots$$

RGPEP

$$t = s^4$$

drop hats

$$H_0(a_0) = \sum_{n=2}^{\infty} \sum_{i_1, i_2, \dots, i_n} c_0(i_1, \dots, i_n) a_{0i_1}^\dagger \cdots a_{0i_n}$$

$$H_t(a_t) = \sum_{n=2}^{\infty} \sum_{i_1, i_2, \dots, i_n} c_t(i_1, \dots, i_n) a_{ti_1}^\dagger \cdots a_{ti_n}$$

$c_0 = \text{initial condition, including } CT_R$

$c_t = ?$

$$H_t(a_t) = H_0(a_0)$$

$$a_t = U_t a_0 U_t^\dagger$$

$$H_t(a_0) = U_t^\dagger H_0(a_0) U_t$$

$$H_t' = \left[-U_t^\dagger U_t', H_t \right] = [G_t, H_t]$$

$$U_t = T \exp \left(- \int_0^t d\tau G_\tau \right)$$

$$G_t = [H_f, \tilde{H}_t]$$

RGPEP generator and non-perturbative QCD

$$G_t = [H_f, \tilde{H}_t]$$

$$H_f = \sum_i p_i^- a_{0i}^\dagger a_{0i} \quad p_i^- = \frac{p_i^{\perp 2}}{p_i^+}$$

$$H_t(a_0) = \sum_{n=2}^{\infty} \sum_{i_1, i_2, \dots, i_n} c_t(i_1, \dots, i_n) a_{0i_1}^\dagger \cdots a_{0i_n}$$

$$\tilde{H}_t(a_0) = \sum_{n=2}^{\infty} \sum_{i_1, i_2, \dots, i_n} c_t(i_1, \dots, i_n) \left(\sum_{k=1}^n p_{i_k}^+ / 2 \right)^2 a_{0i_1}^\dagger \cdots a_{0i_n}$$

$$H'_t = [[H_f, \tilde{H}_t], H_t]$$

$$H_t = H_f + gH_{1t} + g^2H_{2t} + g^3H_{3t} + \dots$$

$$H'_{1t} = [[H_f, \tilde{H}_{1t}], H_f]$$

$$H_{1t} = f_t H_{10}$$

$$f_t = e^{-t(\mathcal{M}_c^2 - \mathcal{M}_a^2)^2}$$

$$H_{A^3 1t} = \sum_{123} \int [123] \delta(1 + 2 - 3) e^{-t \mathcal{M}_{12}^4} \left[Y_{123} a_1^\dagger a_2^\dagger a_3 + Y_{123}^* a_3^\dagger a_2 a_1 \right]$$

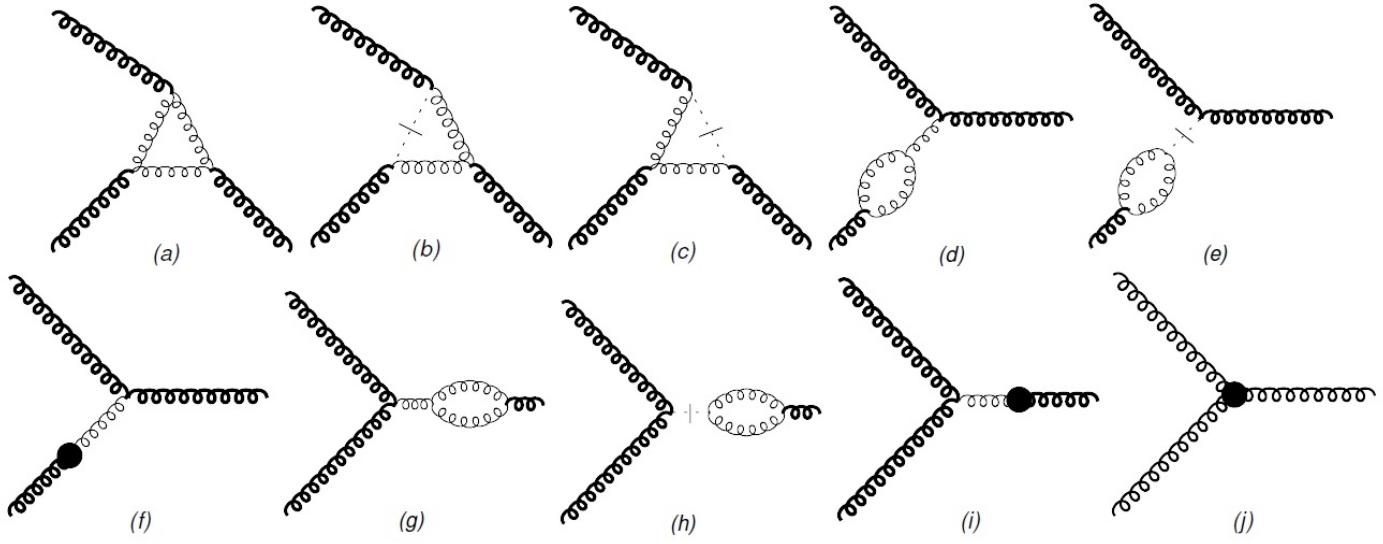


FIG. 1: Third-order contributions to the three-gluon vertex

$$H_{A^3(1+3)t} = \sum_{123} \int [123] \delta_{12.3} e^{-t \mathcal{M}_{12}^4} V_t(x, \kappa^\perp) \left[Y_{123} a_1^\dagger a_2^\dagger a_3 + Y_{123}^* a_3^\dagger a_2 a_1 \right]$$

$$g_t = V_t(x, 0^\perp)$$

$$g_t = g_0 + \frac{g_0^3}{48\pi^2} N_c 11 \ln \frac{s}{s_0} \quad t = s^4$$

Hamiltonian $\beta(s)$ \leftrightarrow Gross-Wilczek-Politzer $\beta(\lambda)$

Minkowski s \leftrightarrow $1/\lambda$ Euclid

Hydrogen atom analogy

$$V_c = -\frac{\alpha_{atom}}{r}$$

$$\psi(\vec{r}) \sim e^{-\alpha_{atom}\mu|\vec{r}|}$$

$$\tilde{\psi}(\vec{k}) \sim \frac{1}{[\vec{k}^2 + \alpha_{atom}^2\mu^2]^2}$$

$$\left[P^+ \hat{H}_{QED}(s_{atom}) - P^{\perp 2} \right] |\psi P\rangle = M_{atom}^2 |\psi P\rangle$$

difference between QED and QCD

RGPEP(g_t^4) \rightarrow hadrons

Conclusion

A. third-order RGPEP is just a beginning (4th order \rightarrow breakthrough?)

B. simple generator \rightarrow non-perturbative definition of QCD

C. two-fold universality:

1) Hamiltonian $\leftrightarrow \overline{MS}$ Minkowski $s \leftrightarrow$ Euclid $1/\lambda$

2) different RGPEP generators

D. SM mass-generation issues \sim QCD mass generation mechanism (?)

E. SM theory issues \sim hierarchy, symmetry breaking, and lc (?)